



First determination of the CP content of $D \rightarrow \pi^+\pi^-\pi^+\pi^-$ and updated determination of the CP contents of $D \rightarrow \pi^+\pi^-\pi^0$ and $D \rightarrow K^+K^-\pi^0$



S. Malde^a, C. Thomas^a, G. Wilkinson^{a,b,*}, P. Naik^c, C. Prouve^c, J. Rademacker^c, J. Libby^d, M. Nayak^d, T. Gershon^e, R.A. Briere^f

^a University of Oxford, Denys Wilkinson Building, Keble Road, OX1 3RH, United Kingdom

^b European Organisation for Nuclear Research (CERN), CH-1211, Geneva 23, Switzerland

^c University of Bristol, Bristol, BS8 1TL, United Kingdom

^d Indian Institute of Technology Madras, Chennai 600036, India

^e University of Warwick, Coventry, CV4 7AL, United Kingdom

^f Carnegie Mellon University, Pittsburgh, PA 15213, USA

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ABSTRACT

Quantum-correlated $\psi(3770) \rightarrow D\bar{D}$ decays collected by the CLEO-c experiment are used to perform a first measurement of $F_+^{4\pi}$, the fractional CP -even content of the self-conjugate decay $D \rightarrow \pi^+\pi^-\pi^+\pi^-$, obtaining a value of 0.737 ± 0.028 . An important input to the measurement comes from the use of $D \rightarrow K_S^0\pi^+\pi^-$ and $D \rightarrow K_L^0\pi^+\pi^-$ decays to tag the signal mode. This same technique is applied to the channels $D \rightarrow \pi^+\pi^-\pi^0$ and $D \rightarrow K^+K^-\pi^0$, yielding $F_+^{\pi\pi\pi^0} = 1.014 \pm 0.045 \pm 0.022$ and $F_+^{KK\pi^0} = 0.734 \pm 0.106 \pm 0.054$, where the first uncertainty is statistical and the second systematic. These measurements are consistent with those of an earlier analysis, based on CP -eigenstate tags, and can be combined to give values of $F_+^{\pi\pi\pi^0} = 0.973 \pm 0.017$ and $F_+^{KK\pi^0} = 0.732 \pm 0.055$. The results will enable the three modes to be included in a model-independent manner in measurements of the unitarity triangle angle γ using $B^\mp \rightarrow DK^\mp$ decays, and in time-dependent studies of CP violation and mixing in the $D^0\bar{D}^0$ system.

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1. Introduction

Studies of the process $B^\mp \rightarrow DK^\mp$, where D indicates a neutral charmed meson reconstructed in a state accessible to both D^0 and \bar{D}^0 decays, give sensitivity to the unitarity triangle angle $\gamma \equiv \arg(-V_{ud}V_{ub}^*/V_{cd}V_{cb}^*)$ (also denoted ϕ_3). Improved knowledge of γ is necessary for testing the Standard Model description of CP violation. In a recent publication [1] it was shown how inclusive three-body self-conjugate D meson decays can be used for this purpose, provided their fractional CP -even content is known, a quantity denoted F_+ (or F_+^f when it is necessary to designate the specific decay f). Measurements of F_+ for the decays $D \rightarrow \pi^+\pi^-\pi^0$ and $D \rightarrow K^+K^-\pi^0$ were performed, making use of quantum-correlated $D\bar{D}$ decays coherently produced at

the $\psi(3770)$ resonance and collected by the CLEO-c detector. In this Letter a first measurement is presented of the CP content of the four-body mode $D \rightarrow \pi^+\pi^-\pi^+\pi^-$, again exploiting CLEO-c $\psi(3770)$ data. This fully-charged and relatively abundant final state [2] can be reconstructed with good efficiency by the LHCb detector and hence is a promising mode for improving the determination of γ at that experiment, as well as at Belle II.

The three-body analysis reported in Ref. [1] exploited events in which one D meson is reconstructed in the signal mode and the other ‘tagging’ meson in its decay to a CP eigenstate. The measurement of $F_+^{4\pi}$ presented in this Letter follows the same method, but augments it with other approaches, in particular a complementary strategy in which the tagging modes are $D \rightarrow K_{S,L}^0\pi^+\pi^-$, and attention is paid to where on the Dalitz plot this tag decay occurs. In order to benefit from this strategy for the previously studied decays, this Letter also presents measurements of $F_+^{\pi\pi\pi^0}$ and $F_+^{KK\pi^0}$ using $D \rightarrow K_{S,L}^0\pi^+\pi^-$ tags. Throughout the effects of CP violation in the charm system are neglected, which is a good

* Corresponding author.

E-mail address: guy.wilkinson@cern.ch (G. Wilkinson).

assumption given theoretical expectations and current experimental limits [2–4]. However, as discussed in Ref. [5], knowledge of F_+ also allows such D decays to be used to study CP -violating observables and mixing parameters through time-dependent measurements at facilities where the mesons are produced incoherently.

The remainder of the Letter is structured as follows. Section 2 introduces the CP -even fraction F_+ , derives the relations that are used to measure its value at the $\psi(3770)$ resonance, and reviews how knowledge of F_+ allows non- CP eigenstates to be cleanly employed in the measurement of γ with $B^\mp \rightarrow DK^\mp$ decays. Section 3 describes the data set and event selection. Sections 4, 5 and 6 presents the determination of F_+ using CP tags, $D \rightarrow K_{S,L}^0 \pi^+ \pi^-$ tags and other tags, respectively. In Section 7 combinations of the individual sets of results are performed for each signal mode; for $D \rightarrow \pi^+ \pi^- \pi^0$ and $D \rightarrow K^+ K^- \pi^0$ these combinations include the results from Ref. [1]. Section 8 gives the conclusions.

2. Measuring the CP content of a self-conjugate D -meson decay and the consequences for the γ determination with $B^\mp \rightarrow DK^\mp$

Let the amplitude of a D^0 meson decaying to a self-conjugate final state f be written as $\mathcal{A}(D^0 \rightarrow f(\mathbf{x})) \equiv a_x e^{i\theta_x}$, where \mathbf{x} indicates a particular point in the decay phase space and θ_x is a CP -conserving strong phase. The amplitude is normalised such that

$$\int_{\mathbf{x} \in \mathcal{D}} |\mathcal{A}(D^0 \rightarrow f(\mathbf{x}))|^2 d\mathbf{x} = \mathcal{B}(f), \quad (1)$$

where $\mathcal{B}(f)$ is the branching fraction of the D^0 decay and \mathcal{D} indicates the entire phase space. The D^0 decay amplitude at $\bar{\mathbf{x}}$ is denoted $a_{\bar{x}} e^{i\theta_{\bar{x}}}$, where $\bar{\mathbf{x}}$ indicates the point in phase space reached by applying a CP transformation to the final-state system at \mathbf{x} . CP violation in the charm system is neglected, which implies that the \bar{D}^0 decay amplitude at $\bar{\mathbf{x}}$ is equal to the D^0 amplitude at \mathbf{x} . It is useful to define the strong phase difference $\Delta\theta_x \equiv \theta_x - \theta_{\bar{x}}$.

It is possible to express the CP -even fraction in terms of the decay amplitudes introduced above. Let the CP eigenstates be $|D_{CP\pm}\rangle \equiv (|D^0\rangle \pm |\bar{D}^0\rangle)/\sqrt{2}$ and consider the decay $D^0 \rightarrow f$ in terms of these states. The total CP -even fraction of the inclusive decay is defined as

$$F_+^f \equiv \frac{\int_{\mathbf{x} \in \mathcal{D}} |\langle f(\mathbf{x}) | D_{CP+} \rangle|^2 d\mathbf{x}}{\int_{\mathbf{x} \in \mathcal{D}} |\langle f(\mathbf{x}) | D_{CP+} \rangle|^2 + |\langle f(\mathbf{x}) | D_{CP-} \rangle|^2 d\mathbf{x}}, \quad (2)$$

and so

$$F_+^f = \frac{\int_{\mathbf{x} \in \mathcal{D}} a_x^2 + a_{\bar{x}}^2 + 2a_x a_{\bar{x}} \cos \Delta\theta_x d\mathbf{x}}{\int_{\mathbf{x} \in \mathcal{D}} 2(a_x^2 + a_{\bar{x}}^2) d\mathbf{x}} = \frac{1}{2} \left[1 + \frac{1}{\mathcal{B}(f)} \int_{\mathbf{x} \in \mathcal{D}} a_x a_{\bar{x}} \cos \Delta\theta_x d\mathbf{x} \right]. \quad (3)$$

Note also that the following relation is always true in the absence of CP violation:

$$\int_{\mathbf{x} \in \mathcal{D}} a_x a_{\bar{x}} \sin \Delta\theta_x d\mathbf{x} = 0. \quad (4)$$

Now consider a quantum-correlated $D\bar{D}$ system produced in the decay of a $\psi(3770)$ meson. One of the D mesons in the system decays to f at the point \mathbf{x} , the other to g at \mathbf{y} , where in general the phase space of the two decays is different. The amplitude of the latter decay is denoted $b_y e^{i\phi_y}$ in analogy with the terminology used above.

The amplitude of the $\psi(3770) \rightarrow D\bar{D} \rightarrow fg$ correlated wavefunction can be written [6]

$$\mathcal{A}(f(\mathbf{x})|g(\mathbf{y})) = \frac{1}{\sqrt{2}} \left[a_x e^{i\theta_x} b_y e^{i\phi_y} - a_{\bar{x}} e^{i\theta_{\bar{x}}} b_y e^{i\phi_y} \right]. \quad (5)$$

The resulting decay probability is then

$$\begin{aligned} \mathcal{P}(f(\mathbf{x})|g(\mathbf{y})) &\propto \left[a_x^2 b_y^2 + a_{\bar{x}}^2 b_y^2 \right. \\ &\quad \left. - 2a_x b_y a_{\bar{x}} b_y \left(\cos \Delta\theta_x \cos \Delta\phi_y + \sin \Delta\theta_x \sin \Delta\phi_y \right) \right]. \end{aligned} \quad (6)$$

If both D mesons decay to the same final state the probability is divided by two to avoid double counting. This formula can be used to determine the population of quantum-correlated decays either integrated over all phase space or after dividing the phase space into bins.

The number of ‘double-tagged’ candidates in which one D meson decays to f and the other to g , integrating over the phase space of each decay, is

$$M(f|g) = \mathcal{Z} \mathcal{B}(f) \mathcal{B}(g) \left[1 - \left(2F_+^f - 1 \right) \left(2F_+^g - 1 \right) \right], \quad (7)$$

where \mathcal{Z} is a normalisation constant common to all decay modes. An important special case, considered in Section 4, is where the tagging-mode g is a CP eigenstate, and $(2F_+^g - 1)$ reduces to ± 1 . Section 6 describes an analysis of classes of double-tags where this is not the case.

Alternatively, when the tagging-mode g is a multibody decay, its phase space may be divided into bins. Integrating over the phase space of f results in the following decay probability in bin i of the phase space of g :

$$\mathcal{P}(f|g_i) \propto \int_{\mathbf{y} \in \mathcal{D}_i} b_y^2 + b_{\bar{y}}^2 - \left(2F_+^f - 1 \right) b_y b_{\bar{y}} \cos \Delta\phi_y d\mathbf{y}, \quad (8)$$

where \mathcal{D}_i indicates the phase space encompassed by bin i . In Section 5 this relation is exploited for the tags $D \rightarrow K_{S,L}^0 \pi^+ \pi^-$.

To understand the relevance of the CP -even fraction in the measurement of the unitarity-triangle angle γ consider the decay of a B^- meson to DK^- , following which the D meson decays to a self-conjugate final state f consisting of three or more particles. The amplitude of the B^- decay is a superposition of two decay paths:

$$\begin{aligned} \mathcal{A}(B^-) &= \mathcal{A}(B^- \rightarrow D^0 K^-) \mathcal{A}(D^0 \rightarrow f) \\ &\quad + \mathcal{A}(B^- \rightarrow \bar{D}^0 K^-) \mathcal{A}(\bar{D}^0 \rightarrow f). \end{aligned} \quad (9)$$

Following the formalism developed above, the decay amplitude of the D^0 meson at the point \mathbf{x} in the phase space is denoted $a_x e^{i\theta_x}$. The decay amplitude of the B^- meson at this point in phase space is therefore

$$\mathcal{A}(B^-(\mathbf{x})) = \mathcal{A}(B^- \rightarrow D^0 K^-) \left[a_x e^{i\theta_x} + r_B e^{i(\delta_B - \gamma)} a_{\bar{x}} e^{i\theta_{\bar{x}}} \right], \quad (10)$$

where r_B and δ_B are respectively the ratio of moduli and the strong phase difference between the suppressed and favoured B^- decay amplitudes. The resulting decay probability is

$$\begin{aligned} \mathcal{P}(B^-(\mathbf{x})) &\propto a_x^2 + r_B^2 a_{\bar{x}}^2 + 2r_B a_x a_{\bar{x}} \cos(\delta_B - \gamma + \theta_x - \theta_{\bar{x}}) \\ &= a_x^2 + r_B^2 a_{\bar{x}}^2 + 2r_B a_x a_{\bar{x}} \left[\cos(\delta_B - \gamma) \cos \Delta\theta_x \right. \\ &\quad \left. - \sin(\delta_B - \gamma) \sin \Delta\theta_x \right]. \end{aligned} \quad (11)$$

Table 1*D*-meson final states reconstructed in this analysis.

Type	Final states
Mixed CP	$\pi^+\pi^-\pi^+\pi^-$, $\pi^+\pi^-\pi^0$, $K^+K^-\pi^0$, $K_S^0\pi^+\pi^-$
CP-even	K^+K^- , $\pi^+\pi^-$, $K_S^0\pi^0\pi^0$, $K_L^0\pi^0$, $K_L^0\omega$
CP-odd	$K_S^0\pi^0$, $K_S^0\omega$, $K_S^0\eta$, $K_S^0\eta'$

The expression for $B^+(\mathbf{x})$ is identical except that the sign in front of γ is reversed and $\mathbf{x} \leftrightarrow \bar{\mathbf{x}}$. The total yield of B^\mp decays is determined by integrating over the entire D phase space:

$$Y^\mp = h^\mp \int_{\mathbf{x} \in \mathcal{D}} \mathcal{P}(B^\mp(\mathbf{x})) d\mathbf{x} = h^\mp \left[1 + r_B^2 + \left(2F_+^f - 1 \right) 2r_B \cos(\delta_B \mp \gamma) \right], \quad (12)$$

where h^\mp is a normalisation constant and Eqs. (3) and (4) have been employed. This expression is very similar to that derived in Ref. [7] for the case when the D meson decays to a CP eigenstate and is indeed identical in the event $F_+^f = 0$ or 1. Hence measurements of Y^\mp , and observables built from these yields [1], can be used to obtain information on the angle γ and the other parameters of the B^\mp decay, provided that F_+^f is known. In Ref. [1] it is demonstrated how the effects of $D^0\bar{D}^0$ mixing, neglected in Eq. (12), may also be accommodated.

3. Data set and event selection

The data set analysed consists of e^+e^- collisions produced by the Cornell Electron Storage Ring (CESR) at $\sqrt{s} = 3.77$ GeV corresponding to an integrated luminosity of 818 pb^{-1} and collected with the CLEO-c detector. The CLEO-c detector is described in detail elsewhere [8]. Monte Carlo simulated samples of signal decays are used to estimate selection efficiencies. Possible background contributions are determined from a generic $D^0\bar{D}^0$ simulated sample corresponding to approximately fifteen times the integrated luminosity of the data set. The EVTGEN generator [9] is used to simulate the decays. The detector response is modelled using the GEANT software package [10].

Table 1 lists the D -meson final states considered in the analysis. Double-tag candidates are reconstructed in which one D meson decays into $\pi^+\pi^-\pi^+\pi^-$ and the other into a CP eigenstate, or where one D meson decays into $\pi^+\pi^-\pi^+\pi^-$, $\pi^+\pi^-\pi^0$ or $K^+K^-\pi^0$ and the other into one of the mixed-CP modes $K_S^0\pi^+\pi^-$ or $K_L^0\pi^+\pi^-$. The combinations $\pi^+\pi^-\pi^+\pi^-$ vs. $\pi^+\pi^-\pi^+\pi^-$ and $\pi^+\pi^-\pi^+\pi^-$ vs. $\pi^+\pi^-\pi^0$ are also reconstructed.

The unstable final state particles are reconstructed in the following decay modes: $\pi^0 \rightarrow \gamma\gamma$, $K_S^0 \rightarrow \pi^+\pi^-$, $\omega \rightarrow \pi^+\pi^-\pi^0$, $\eta \rightarrow \gamma\gamma$, $\eta \rightarrow \pi^+\pi^-\pi^0$ and $\eta' \rightarrow \eta(\gamma\gamma)\pi^+\pi^-$. The π^0 , K_S^0 , ω ,

η and η' reconstruction procedure is identical to that used in Ref. [11].

Final states that do not contain a K_L^0 are fully reconstructed via the beam-constrained candidate mass, $m_{bc} \equiv \sqrt{s/(4c^4) - \mathbf{p}_D^2/c^2}$, where \mathbf{p}_D is the D -candidate momentum, and $\Delta E \equiv E_D - \sqrt{s}/2$, where E_D is the D -candidate energy. The m_{bc} and ΔE distributions of correctly reconstructed D -meson candidates peak at the nominal D^0 mass and zero, respectively. Neither ΔE nor m_{bc} distributions exhibit any peaking structure for combinatoric background. The double-tag yield is determined from counting events in signal and sideband regions of m_{bc} after placing requirements on ΔE [1,11–13]. The selection criteria of candidates involving the modes $D \rightarrow K^+K^-$ and $D \rightarrow \pi^+\pi^-$ do not include the cosmic ray muon and radiative Bhabha vetoes that are described in Ref. [1]. This is because these sources of background do not contaminate the double-tag sample, and the vetoes are found to perturb the selection efficiency of the other D meson in the event. When selecting $D \rightarrow K_S^0\pi^+\pi^-$ candidates it is demanded that the K_S^0 decay products form a vertex that is significantly displaced from the e^+e^- collision point; in contrast, for $D \rightarrow \pi^+\pi^-\pi^+\pi^-$ and $D \rightarrow \pi^+\pi^-\pi^0$ candidates the $\pi^+\pi^-$ vertex must be consistent with originating from the collision point in order to suppress contamination from $D \rightarrow K_S^0\pi^+\pi^-$ and $D \rightarrow K_S^0\pi^0$ decays, respectively.

The double-tag yield determination procedure is identical to that presented in Refs. [11,12] except for the selections where the signal decay is $\pi^+\pi^-\pi^+\pi^-$ and the tag decay is K^+K^- , $\pi^+\pi^-$, $\pi^+\pi^-\pi^0$ or $\pi^+\pi^-\pi^+\pi^-$, which are all dominated by a background from continuum production of light quark-antiquark pairs. For these modes an unbinned maximum likelihood fit is performed to the distribution of the average m_{bc} of the two D -meson candidates. The background is modelled with an ARGUS function [14] and the signal is modelled with the sum of two Crystal Ball functions [15] with power-law tails on opposite sides. The parameters of the Crystal Ball functions are fixed from fits to large samples of simulated data.

Figs. 1(a) and (b) show the average m_{bc} distributions for CP-tagged $D \rightarrow \pi^+\pi^-\pi^+\pi^-$ candidates, summed over all tag modes that are CP-even and CP-odd eigenstates, respectively, where the CP-tag final state does not contain a K_L^0 meson. Fig. 2 shows the average m_{bc} distributions for $D \rightarrow \pi^+\pi^-\pi^+\pi^-$, $D \rightarrow \pi^+\pi^-\pi^0$ and $D \rightarrow K^+K^-\pi^0$ candidates tagged with $D \rightarrow K_S^0\pi^+\pi^-$ decays, while Figs. 3(a)–(c) show the Dalitz-plot distributions of the tag decay for these three signal modes.

Many K_L^0 mesons do not deposit any reconstructible signal in the detector. However, double-tag candidates can be fully reconstructed using a missing-mass squared (m_{miss}^2) technique [16] for tags containing a single K_L^0 meson. Yields are determined from the signal and sideband regions of the m_{miss}^2 distribution. Fig. 1(c) shows the m_{miss}^2 distributions for $D \rightarrow \pi^+\pi^-\pi^+\pi^-$ can-

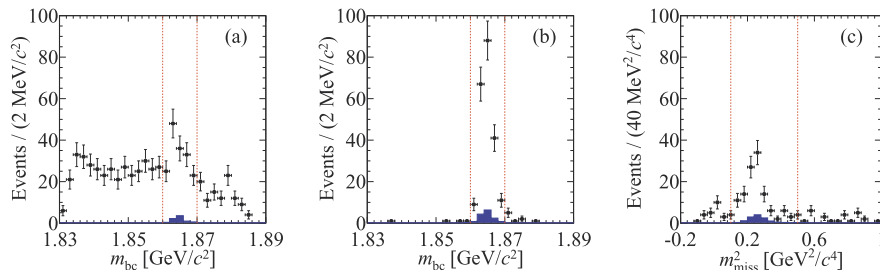


Fig. 1. Distributions of $D \rightarrow \pi^+\pi^-\pi^+\pi^-$ candidates tagged by CP-eigenstates. Sub-figures (a) and (b) show average m_{bc} distributions for CP-even tags and CP-odd tags not involving K_L^0 mesons, respectively. Sub-figure (c) shows the m_{miss}^2 distribution for candidates tagged by CP eigenstates that contain a K_L^0 meson. The shaded histogram is the estimated peaking background and the vertical dotted lines indicate the signal region.

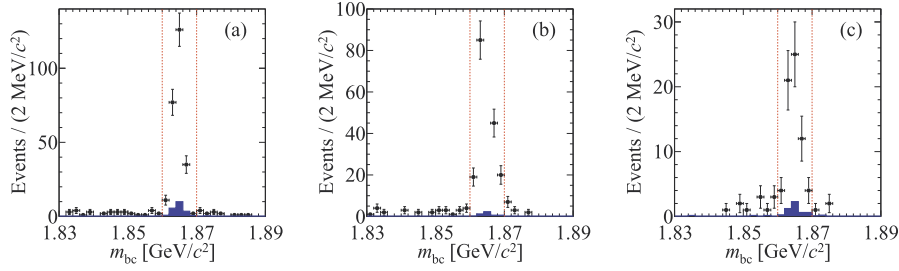


Fig. 2. Average m_{bc} distributions for (a) $D \rightarrow \pi^+\pi^-\pi^+\pi^-$, (b) $D \rightarrow \pi^+\pi^-\pi^0$ and (c) $D \rightarrow K^+K^-\pi^0$ candidates tagged by a $D \rightarrow K_S^0\pi^+\pi^-$ decay. The shaded histogram is the estimated peaking background and the vertical dotted lines indicate the signal region.

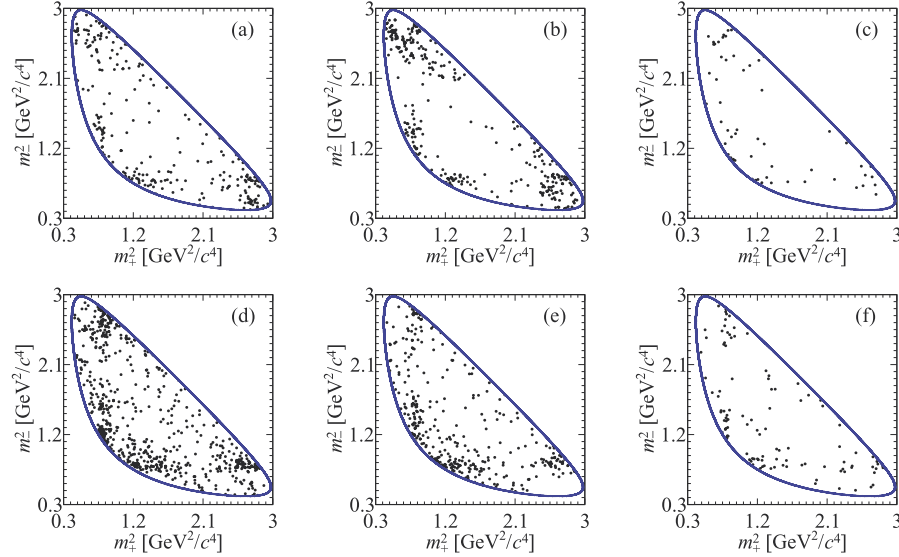


Fig. 3. Dalitz-plot distributions for $D \rightarrow K_S^0\pi^+\pi^-$ reconstructed against (a) $D \rightarrow \pi^+\pi^-\pi^+\pi^-$, (b) $D \rightarrow \pi^+\pi^-\pi^0$ and (c) $D \rightarrow K^+K^-\pi^0$, and $D \rightarrow K_L^0\pi^+\pi^-$ reconstructed against (d) $D \rightarrow \pi^+\pi^-\pi^+\pi^-$, (e) $D \rightarrow \pi^+\pi^-\pi^0$ and (f) $D \rightarrow K^+K^-\pi^0$. The axis labels m_{\pm}^2 are the invariant-mass squared of the $\pi^{\pm}K_{S,L}^0$ pair.

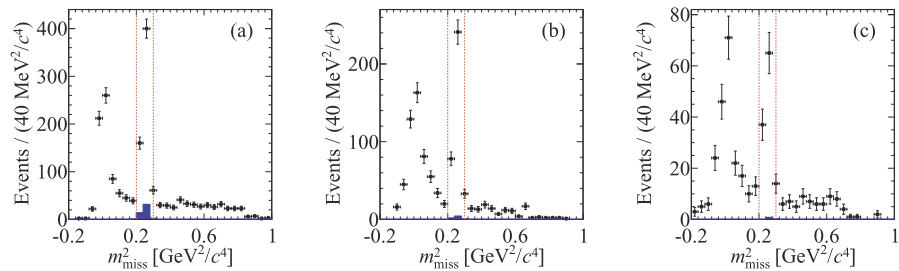


Fig. 4. m_{miss}^2 distributions for (a) $D \rightarrow \pi^+\pi^-\pi^+\pi^-$, (b) $D \rightarrow \pi^+\pi^-\pi^0$ and (c) $D \rightarrow K^+K^-\pi^0$ candidates tagged by a $D \rightarrow K_L^0\pi^+\pi^-$ decay. The shaded histogram is the estimated peaking background and the vertical dotted lines indicate the signal region.

didates tagged with either a $K_L^0\pi^0$ or $K_L^0\omega$ decay. Fig. 4 shows the m_{miss}^2 distributions for $D \rightarrow \pi^+\pi^-\pi^+\pi^-$, $D \rightarrow \pi^+\pi^-\pi^0$ and $D \rightarrow K^+K^-\pi^0$ candidates tagged with $D \rightarrow K_L^0\pi^+\pi^-$ decays, and Figs. 3(d)–(f) show the corresponding tag-side Dalitz-plot distributions.

In events where more than one pair of decays is reconstructed an algorithm is applied to select a single double-tag candidate based on the information provided by the m_{bc} and ΔE variables. The particular choice of metric varies depending on the category of double tag and is optimised through simulation studies.

The peaking background estimates are determined from the generic Monte Carlo sample of $D^0\bar{D}^0$ events. For double tags involving a CP mode without a K_L^0 meson the peaking backgrounds are found to constitute 5–10% of the selected events, and are pre-

dominantly from residual $D \rightarrow K_S^0\pi^+\pi^-$ contamination. The peaking backgrounds for final states with a K_L^0 are generally larger; for $K_L^0\pi^0$ and $K_L^0\omega$ this contamination amounts to 15–20% of the signal yield, whereas for $K_L^0\pi^+\pi^-$ it is $\sim 10\%$ of the signal yield. The dominant source of peaking background in each case is the equivalent decay containing a K_S^0 instead of a K_L^0 meson. The contamination from specific modes in the other categories of double tags is typically 10% or less. The statistical uncertainties on these background estimates arising from the finite size of the simulated samples are included in the total statistical uncertainties on the signal yields.

The measured double-tag event yields after background subtraction are given in Table 2.

Table 2

Double-tagged signal yields after background subtraction. Information on the entries marked †, not studied in the current analysis, can be found in Ref. [1].

	$\pi^+\pi^-\pi^+\pi^-$	$\pi^+\pi^-\pi^0$	$K^+K^-\pi^0$
K^+K^-	19.3 ± 6.3	†	†
$\pi^+\pi^-$	3.3 ± 8.2	†	†
$K_S^0\pi^0\pi^0$	18.6 ± 5.2	†	†
$K_L^0\pi^0$	49.2 ± 10.9	†	†
$K_L^0\omega$	22.0 ± 6.5	†	†
$K_S^0\pi^0$	112.8 ± 11.0	†	†
$K_S^0\omega$	41.0 ± 6.8	†	†
$K_S^0\eta(\gamma\gamma)$	18.8 ± 4.5	†	†
$K_S^0\eta(\pi^+\pi^-\pi^0)$	3.1 ± 2.7	†	†
$K_S^0\eta'$	9.3 ± 3.3	†	†
$K_S^0\pi^+\pi^-$	217.9 ± 16.8	289.2 ± 17.6	52.5 ± 7.8
$K_L^0\pi^+\pi^-$	485.0 ± 26.3	291.1 ± 19.2	78.1 ± 11.1
$\pi^+\pi^-\pi^+\pi^-$	41.0 ± 16.3	75.5 ± 15.7	-

Knowledge of the single-tag yields of the CP -eigenstate modes is required for normalisation purposes. Since the single-tag reconstruction criteria applied are identical to those employed in Ref. [1], all information on these yields is taken from the earlier publication. It is also necessary to know the single-tag yield for the decay $D \rightarrow \pi^+\pi^-\pi^0$. A fit to the m_{bc} distribution returns a result of 29998 ± 320 signal candidates, after the subtraction of small peaking-background contributions.

4. Analysis with the CP tags

The yields of the single and double tags containing a CP eigenstate are used as inputs to determine the CP -even fraction, $F_+^{4\pi}$. Following on from Eq. (7), the expected number of observed events, M , where one D meson decays to the $\pi^+\pi^-\pi^+\pi^-$ final state, and the other decays to X , a CP eigenstate with eigenvalue η_{CP} , is given by

$$M(4\pi|X) = 2N_{D\bar{D}} \mathcal{B}(4\pi)\mathcal{B}(X)\varepsilon(4\pi|X) \left[1 - \eta_{CP}(2F_+^{4\pi} - 1) \right], \quad (13)$$

where $N_{D\bar{D}}$ is the number of $D\bar{D}$ pairs, $\mathcal{B}(4\pi)$ and $\mathcal{B}(X)$ are the branching fractions for the two reconstructed final states and $\varepsilon(4\pi|X)$ is the efficiency of reconstructing such a double tag. The double tag yield is denoted by M^- (M^+) for CP -even (CP -odd) tags. Experimentally it is advantageous to eliminate dependence on $N_{D\bar{D}}$, the branching fractions and the reconstruction efficiency, which can be achieved by normalising by the single-tag yields. The

yield of single tags, S^+ (S^-) decaying to a CP -odd (CP -even) eigenstate X , is given by

$$S(X) = 2N_{D\bar{D}} \mathcal{B}(X)\varepsilon(X), \quad (14)$$

where $\varepsilon(X)$ is the reconstruction efficiency of the single tag. The small effects of $D^0\bar{D}^0$ mixing are eliminated from the measurement by correcting the measured single-tag yields S_{meas}^\pm such that $S^\pm = S_{\text{meas}}^\pm / (1 - \eta_{CP}y_D)$ where $y_D = (0.62 \pm 0.08)\%$ is the well-known D -mixing parameter [17]. A further correction is applied in the case of the tags K^+K^- and $\pi^+\pi^-$ because of the differing selection requirements for the single and double-tag case, as described in Section 3. This correction factor is determined by taking the ratio of the selection efficiency of the single tag from simulation with the two differing selections. It is determined to be 1.15 and 1.10 for the $D \rightarrow K^+K^-$ and $D \rightarrow \pi^+\pi^-$ modes, respectively, with an uncertainty of ± 0.05 . The other uncertainties on these single-tag yields are assigned following the same procedure described in Ref. [1].

For the case of the two CP tags involving a K_L^0 meson a different treatment is required, since it is not possible to measure the single-tag yield directly for these modes. Following the procedure described in Ref. [1], the effective single-tag yield is evaluated using Eq. (14), where the effective single-tag efficiency $\varepsilon(K_L^0X)$ is calculated from the ratio of $\varepsilon(4\pi|K_L^0X)/\varepsilon(4\pi)$, and the leading systematic uncertainties are associated with the branching fractions and the value used for the effective single-tag efficiency. The effective single-tag yields are determined to be 21726 ± 3497 and 9124 ± 4105 for $K_L^0\pi^0$ and $K_L^0\omega$, respectively.

Assuming that the reconstruction efficiencies of each D meson are independent, then the ratio of the double-tagged and single-tagged yields are independent of the branching fraction and reconstruction efficiency of the CP tag and $N_{D\bar{D}}$. This ratio is defined as $N^+ \equiv M^+/S^+$, with an analogous expression for N^- . The CP -even fraction $F_+^{4\pi}$ is then given by

$$F_+^{4\pi} = \frac{N^+}{N^+ + N^-}. \quad (15)$$

The measured values for N^+ and N^- for each CP tag are displayed in Fig. 5. It can be seen that there is consistency between the individual tags for each measurement. The mean value $\langle N^+ \rangle = (5.54 \pm 0.46) \times 10^{-3}$ is significantly larger than $\langle N^- \rangle = (1.80 \pm 0.32) \times 10^{-3}$, indicating that the $\pi^+\pi^-\pi^+\pi^-$ final state is predominantly CP even.

If the acceptance across the phase space of the $D \rightarrow \pi^+\pi^-\pi^+\pi^-$ decay is not uniform it has the potential to bias the measurement of $F_+^{4\pi}$. Using simulated data the selection ef-

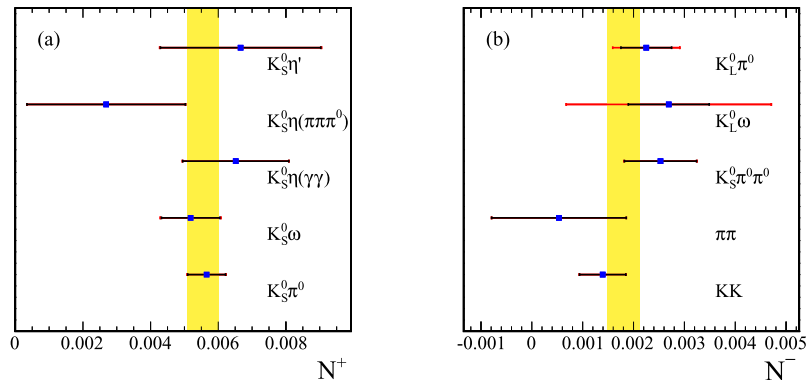


Fig. 5. $D \rightarrow \pi^+\pi^-\pi^+\pi^-$ results for (a) N^+ and (b) N^- . In each plot the vertical yellow band indicates the value obtained from the combination of all tags. The black portion of the uncertainty represents the statistical uncertainty only while the red represents the total. (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)

efficiency of individual pions in $D \rightarrow \pi^+\pi^-\pi^+\pi^-$ decays is determined in bins of momentum and polar angle with respect to the beam direction. The candidates in data are then weighted by the normalised efficiency. Each pion is treated independently and an overall weight, typically lying within 5–10% of unity, is found by multiplying the individual weights. The scaled signal yields are used to re-determine $F_+^{4\pi}$ and the difference between this and the value found without efficiency correction is 0.008, which is taken as the systematic uncertainty due to non-uniform acceptance.

Using the CP tags only, and accounting for the correlations between the systematic uncertainties, yields $F_+^{4\pi} = 0.754 \pm 0.031 \pm 0.021$, where the first uncertainty is statistical and the second systematic.

5. Analysis with the $D \rightarrow K_{S,L}^0\pi^+\pi^-$ tags

For each of the signal samples that are tagged by $D \rightarrow K_S^0\pi^+\pi^-$ or $D \rightarrow K_L^0\pi^+\pi^-$ decays the Dalitz plot of the tag mode is divided into eight pairs of symmetric bins by the line $m_+^2 = m_-^2$, where m_\pm^2 is the invariant-mass squared of the $\pi^\pm K_{S,L}^0$ pair. The bins lying on one side of this line ($m_+^2 > m_-^2$) are labelled $-1 \rightarrow -8$, and those on the other side $1 \rightarrow 8$. The binning definition follows the ‘Equal $\Delta\delta_D$ BABAR 2008’ scheme of Ref. [18], in which the boundaries are chosen according to the strong-phase prediction of a model developed by the BaBar collaboration [19]. The expected distribution of entries is symmetric and so the analysis considers the absolute bin number $|i|$, which contains the contents of the pair of bins $-i$ and i .

Following Eq. (8), the expected population of bin $|i|$ for signal decays with $K_S^0\pi^+\pi^-$ tags is

$$M_{|i|} = h \left[K_i + K_{-i} - (2F_+ - 1)2c_i\sqrt{K_i K_{-i}} \right], \quad (16)$$

where h is a normalisation factor specific to the signal category, K_i is the flavour-tagged fraction, being the proportion of $K_S^0\pi^+\pi^-$ decays to fall in bin i in the case that the mother particle is known to be a D^0 meson, and c_i is the cosine of the strong-phase difference between D^0 and \bar{D}^0 decays averaged in bin i and weighted by the absolute decay rate (a precise definition may be found in Ref. [6]). The only difference between the form of this expression and the case when the signal decays into a pure CP -even eigenstate [6] is the prefactor of $(2F_+ - 1)$ in the final term.

Similarly, when the tagging meson decays to $K_L^0\pi^+\pi^-$ then the number of double-tag decays produced in bin $|i|$ is

$$M'_{|i|} = h' \left[K'_i + K'_{-i} + (2F_+ - 1)2c'_i\sqrt{K'_i K'_{-i}} \right], \quad (17)$$

where the primed quantities are now specific to this case. The reversed sign in front of the final term reflects the fact that the K_L^0 meson is almost entirely a CP -odd eigenstate.

The values of c_i and c'_i within these bins have been measured by the CLEO collaboration [18]. The values of the K_i parameters are taken from an analysis of the predictions of various B -factory models [19–22] presented in Ref. [13], and those of the K'_i parameters from measurements performed with CLEO-c data [23].

The double-tagged samples are analysed to determine the background-subtracted signal yield in each Dalitz-plot bin. The distribution of background between the different bins is assigned according to its category. Flat background is assumed to contribute proportionally to the bin area. Peaking backgrounds that occur on the signal side affect the distribution of tag decays in $K_{S,L}^0\pi^+\pi^-$ phase space according to their nature. For example, in the case of $D \rightarrow K_S^0\pi^0$ decays that are wrongly reconstructed as $D \rightarrow \pi^+\pi^-\pi^0$, the tag decay will be in a CP -even state and distributed accordingly. Similarly, the distribution of $K_S^0(\pi^0\pi^0)\pi^+\pi^-$

Table 3

Double-tagged signal yields vs. $K_S^0\pi^+\pi^-$ after background subtraction in absolute bin numbers of the $D \rightarrow K_S^0\pi^+\pi^-$ Dalitz plot. The yields are corrected for relative bin-to-bin efficiency variations and then scaled so that the totals match the values in Table 2.

$ i $	$\pi^+\pi^-\pi^+\pi^-$	$\pi^+\pi^-\pi^0$	$K^+K^-\pi^0$
1	30.8 ± 7.0	29.9 ± 6.3	12.6 ± 4.1
2	19.8 ± 5.3	19.1 ± 4.8	4.6 ± 2.6
3	16.4 ± 4.5	27.2 ± 5.2	6.9 ± 2.5
4	10.1 ± 3.4	18.5 ± 4.4	1.6 ± 1.5
5	55.1 ± 8.1	96.9 ± 10.0	8.4 ± 3.1
6	21.1 ± 5.1	31.2 ± 5.8	4.4 ± 2.4
7	27.7 ± 6.0	34.6 ± 6.3	7.7 ± 2.8
8	36.9 ± 6.8	31.8 ± 6.0	6.2 ± 2.5

Table 4

Double-tagged signal yields vs. $K_L^0\pi^+\pi^-$ after background subtraction in absolute bin numbers of the $D^0 \rightarrow K_L^0\pi^+\pi^-$ Dalitz plot. The yields are corrected for relative bin-to-bin efficiency variations.

$ i $	$\pi^+\pi^-\pi^+\pi^-$	$\pi^+\pi^-\pi^0$	$K^+K^-\pi^0$
1	134.1 ± 13.9	89.2 ± 11.1	17.3 ± 6.1
2	59.2 ± 8.9	32.9 ± 6.9	8.8 ± 4.0
3	55.4 ± 8.7	31.0 ± 6.3	4.1 ± 3.1
4	20.3 ± 5.8	7.0 ± 3.7	0.1 ± 1.9
5	46.0 ± 8.7	6.7 ± 4.8	2.1 ± 3.1
6	24.6 ± 6.2	14.7 ± 5.0	10.0 ± 3.7
7	61.2 ± 9.0	46.7 ± 7.8	17.6 ± 4.7
8	84.1 ± 10.8	62.9 ± 8.9	18.1 ± 5.1

decays that are misreconstructed as $K_L^0\pi^+\pi^-$ tags is well understood and modelled appropriately. The distribution of the residual $K_S^0\pi^+\pi^-$ vs. $K_S^0\pi^+\pi^-$ events that contaminate the $\pi^+\pi^-\pi^+\pi^-$ vs. $K_S^0\pi^+\pi^-$ selection is determined from data by inverting the K_S^0 veto on the signal decay.

It is also necessary to account for relative bin-to-bin efficiency variations in the background-subtracted signal yields. The correction factors are determined from simulation and typically differ $\lesssim 5\%$ from unity. The signal yields in each bin after background subtraction and relative efficiency correction are shown in Table 3 for $K_S^0\pi^+\pi^-$ tags and in Table 4 for $K_L^0\pi^+\pi^-$ tags.

A log-likelihood fit is performed to the efficiency-corrected signal yields of each sample, assuming the expected distributions given by Eqs. (16) and (17). The fit parameters are the CP -even fraction and the overall normalisation. The values of K_i , K'_i , c_i and c'_i are also fitted, but with their measurement uncertainties and correlations imposed with Gaussian constraints. Separate fits are performed for the $D \rightarrow K_S^0\pi^+\pi^-$ tags, the $D \rightarrow K_L^0\pi^+\pi^-$ tags, and for both samples combined. Fits to large ensembles of simulated experiments demonstrate that the returned uncertainties are reliable and that there is no significant bias in the procedure. All data fits are found to be of good quality. The results are plotted in Fig. 6 for the $D \rightarrow K_S^0\pi^+\pi^-$ tags and in Fig. 7 for the $D \rightarrow K_L^0\pi^+\pi^-$ tags. The numerical results for the CP -even fraction are given in Table 5 for $D \rightarrow \pi^+\pi^-\pi^+\pi^-$ and in Table 6 for $D \rightarrow \pi^+\pi^-\pi^0$ and $D \rightarrow K^+K^-\pi^0$.

The dominant systematic uncertainty is associated with the distribution of the continuum and combinatoric backgrounds. This is assessed by repeating the measurement with these contributions switched to those found in the sidebands of the signal distributions. Shifts of 0.020, 0.019 and 0.037 are observed for $F_+^{4\pi}$, $F_+^{\pi\pi\pi^0}$ and $F_+^{KK\pi^0}$, respectively. The uncertainty associated with the measurement errors on K_i , K'_i , c_i and c'_i is estimated by re-running the fit with these quantities set as fixed parameters and subtracting in quadrature the new fit uncertainty from that obtained with the original procedure. This component is found to be 0.013 for $D \rightarrow \pi^+\pi^-\pi^+\pi^-$, 0.010 for $D \rightarrow \pi^+\pi^-\pi^0$ and

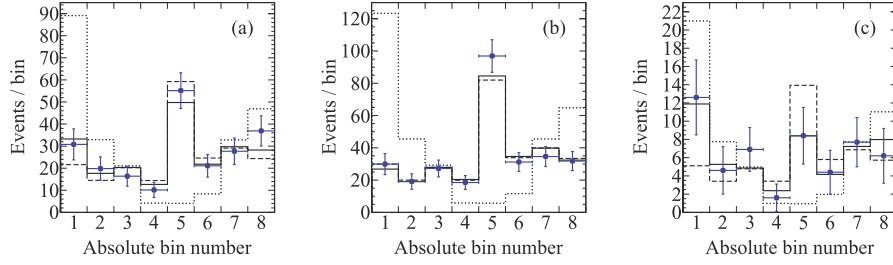


Fig. 6. Data (points) and fit results (solid line) in absolute bin numbers for $K_S^0 \pi^+ \pi^-$ tags vs. (a) $D \rightarrow \pi^+ \pi^- \pi^+ \pi^-$, (b) $D \rightarrow \pi^+ \pi^- \pi^0$ and (c) $D \rightarrow K^+ K^- \pi^0$. Also shown in each case is the expectation if $F_+ = 0$ (dotted line) or $F_+ = 1$ (dashed line).

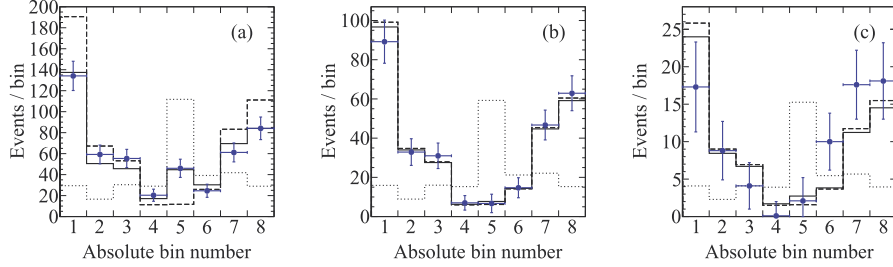


Fig. 7. Data (points) and fit results (solid line) in absolute bin numbers for $K_L^0 \pi^+ \pi^-$ tags vs. (a) $D \rightarrow \pi^+ \pi^- \pi^+ \pi^-$, (b) $D \rightarrow \pi^+ \pi^- \pi^0$ and (c) $D \rightarrow K^+ K^- \pi^0$. Also shown in each case is the expectation if $F_+ = 0$ (dotted line) or $F_+ = 1$ (dashed line).

Table 5

The $F_+^{4\pi}$ fit results for the $D \rightarrow K^0 \pi^+ \pi^-$ tags, where the first uncertainty is statistical and the second systematic. The row $K_{S,L}^0 \pi^+ \pi^-$ indicates the configuration where the CP -even fraction is a common fit parameter shared between the $D \rightarrow K_S^0 \pi^+ \pi^-$ and $D \rightarrow K_L^0 \pi^+ \pi^-$ samples.

Tag	$F_+^{4\pi}$
$K_S^0 \pi^+ \pi^-$	$0.828 \pm 0.074 \pm 0.014$
$K_L^0 \pi^+ \pi^-$	$0.670 \pm 0.057 \pm 0.039$
$K_{S,L}^0 \pi^+ \pi^-$	$0.737 \pm 0.049 \pm 0.024$

Table 6

The $F_+^{\pi\pi\pi^0}$ and $F_+^{KK\pi^0}$ fit results for the $D \rightarrow K^0 \pi^+ \pi^-$ tags, where the first uncertainty is statistical and the second systematic. The row $K_{S,L}^0 \pi^+ \pi^-$ indicates the configuration where the CP -even fraction is a common fit parameter shared between the $D \rightarrow K_S^0 \pi^+ \pi^-$ and $D \rightarrow K_L^0 \pi^+ \pi^-$ samples.

Tag	$F_+^{\pi\pi\pi^0}$	$F_+^{KK\pi^0}$
$K_S^0 \pi^+ \pi^-$	$1.034 \pm 0.054 \pm 0.023$	$0.573 \pm 0.152 \pm 0.046$
$K_L^0 \pi^+ \pi^-$	$0.971 \pm 0.075 \pm 0.033$	$0.916 \pm 0.181 \pm 0.066$
$K_{S,L}^0 \pi^+ \pi^-$	$1.014 \pm 0.045 \pm 0.022$	$0.734 \pm 0.106 \pm 0.054$

0.025 for $D \rightarrow K^+ K^- \pi^0$, and is accounted as a systematic uncertainty in the final results. An uncertainty is evaluated to account for non-uniformities in acceptance across phase space. For $D \rightarrow \pi^+ \pi^- \pi^+ \pi^-$ this contribution is calculated with the same procedure as in Section 4, and found to be 0.002 for the joint $K_{S,L}^0 \pi^+ \pi^-$ fit. For $D \rightarrow \pi^+ \pi^- \pi^0$ and $D \rightarrow K^+ K^- \pi^0$ the acceptance uncertainties are taken to be 0.001 and 0.010, respectively, as determined in Ref. [1]. Other sources of bias are evaluated to be small. The total systematic uncertainties are given in Tables 5 and 6 and are in all cases significantly smaller than the corresponding statistical uncertainties.

6. Other tags

The double-tagged yield of $\pi^+ \pi^- \pi^+ \pi^-$ vs. $\pi^+ \pi^- \pi^0$ can be used to determine $F_+^{4\pi}$, benefiting from the well-measured value of $F_+^{\pi\pi\pi^0}$. The ratio of double-tag and $D \rightarrow \pi^+ \pi^- \pi^0$ single-tag

yields is defined as $N^{\pi\pi\pi^0} \equiv M(4\pi|\pi\pi\pi^0)/S(\pi\pi\pi^0)$, where a very small correction is applied to the measured single-tag yield to account for mixing effects. Following from Eq. (7), the ratio $N^{\pi\pi\pi^0}/N^+$ removes dependence on the signal branching fraction and reconstruction efficiency and is given by

$$\frac{N^{\pi\pi\pi^0}}{N^+} = \frac{[1 - (2F_+^{4\pi} - 1)(2F_+^{\pi\pi\pi^0} - 1)]}{2F_+^{4\pi}}, \quad (18)$$

which can be rearranged to yield

$$F_+^{4\pi} = \frac{N^+ F_+^{\pi\pi\pi^0}}{N^{\pi\pi\pi^0} - N^+ + 2N^+ F_+^{\pi\pi\pi^0}}. \quad (19)$$

The choice of N^+ in the denominator of Eq. (18) is preferred to N^- as it is measured with better relative precision.

Taking as input the yields given in Section 3, the value of N^+ reported in Section 4 and the final result for $F_+^{\pi\pi\pi^0}$ presented in Section 7 implies $F_+^{4\pi} = 0.695 \pm 0.050 \pm 0.021$, where the uncertainties are statistical and systematic, respectively. The main contributions to the systematic uncertainty arise from: the measurement of the $D \rightarrow \pi\pi\pi^0$ single-tag yield and small violations of the efficiency-factorisation ansatz assumed in Eq. (18); the understanding of the peaking background component in the sample; and the possible effects of non-uniform acceptance.

The self-tagged yield of $\pi^+ \pi^- \pi^+ \pi^-$ vs. $\pi^+ \pi^- \pi^+ \pi^-$ also carries information on the value of $F_+^{4\pi}$. This sample is however only used for a consistency check, as there are large backgrounds from both the continuum and from misidentification of $D \rightarrow K_S^0 \pi\pi$ decays that are a potential source of significant systematic bias. Furthermore, the predicted yield and measurement uncertainty means that the result from analysis of these double tags would have low weight in the combined measurement of $F_+^{4\pi}$. Using Eq. (7) the number of observed self-tagged events is given by

$$M(4\pi|4\pi) = 4\mathcal{R}F_+^{4\pi}(1 - F_+^{4\pi}), \quad (20)$$

Table 7

Results for $F_+^{4\pi}$ for each tag category, and combined. When two uncertainties are shown, the first is statistical and the second systematic. For the combined result the total uncertainty is given.

Tag	$F_+^{4\pi}$
CP eigenstates	$0.754 \pm 0.031 \pm 0.021$
$K_{S,L}^0 \pi^+ \pi^-$	$0.737 \pm 0.049 \pm 0.024$
$\pi^+ \pi^- \pi^0$	$0.695 \pm 0.050 \pm 0.021$
Combined	0.737 ± 0.028

Table 8

Results for $F_+^{\pi\pi\pi^0}$ and $F_+^{KK\pi^0}$ for each tag category, and combined. The CP-eigenstate tag results are from Ref. [1]. When two uncertainties are shown, the first is statistical and the second systematic. For the combined result the total uncertainty is given.

Tag	$F_+^{\pi\pi\pi^0}$	$F_+^{KK\pi^0}$
CP eigenstates	$0.968 \pm 0.017 \pm 0.006$	$0.731 \pm 0.058 \pm 0.021$
$K_{S,L}^0 \pi^+ \pi^-$	$1.014 \pm 0.045 \pm 0.022$	$0.734 \pm 0.106 \pm 0.054$
Combined	0.973 ± 0.017	0.732 ± 0.055

where $\mathcal{R} = N_{D\bar{D}} \mathcal{B}(4\pi)^2 \varepsilon(4\pi|4\pi)$. The predicted double-tagged yield using the value of $F_+^{4\pi}$ obtained from the CP tags is 17 ± 2 , which is consistent with the measured yield reported in Table 2.

7. Combination of results

The results for $F_+^{4\pi}$ from the CP tags, the $K_{S,L}^0 \pi^+ \pi^-$ tags and the $\pi^+ \pi^- \pi^0$ tag are summarised in Table 7. They are compatible and are therefore combined, taking account of correlated uncertainties. Correlations arise from the non-flat Dalitz plot acceptance between all three measurements and the use of N^+ as an input to both the CP tags and $\pi^+ \pi^- \pi^0$ tag measurements. There is a further small correlation between the results obtained with the CP and $\pi^+ \pi^- \pi^0$ tags, associated with the uncertainty on the value of the mixing parameter y_D . The final result is $F_+^{4\pi} = 0.737 \pm 0.028$.

Table 8 summarises the results on $F_+^{\pi\pi\pi^0}$ and $F_+^{KK\pi^0}$ obtained with $K_{S,L}^0 \pi^+ \pi^-$ tags, together with those determined from CP tags. The $K_{S,L}^0 \pi^+ \pi^-$ measurements confirm the results of the earlier analysis. A combination is performed, accounting for the sole source of correlated uncertainties, which is that arising from the non-uniform acceptance over the Dalitz plots. Results of $F_+^{\pi\pi\pi^0} = 0.973 \pm 0.017$ and $F_+^{KK\pi^0} = 0.732 \pm 0.055$ are obtained. The $K_{S,L}^0 \pi^+ \pi^-$ tags improve the relative precision on $F_+^{\pi\pi\pi^0}$ by 6% and on $F_+^{KK\pi^0}$ by 10%.

8. Conclusions

A first measurement has been made of the CP-even fraction of the decay $D \rightarrow \pi^+ \pi^- \pi^+ \pi^-$, exploiting quantum-correlated double-tags involving CP-eigenstates, a binned Dalitz-plot analysis of the modes $D \rightarrow K_{S,L}^0 \pi^+ \pi^-$, and $D \rightarrow \pi^+ \pi^- \pi^0$ decays. The result, $F_+^{4\pi} = 0.737 \pm 0.028$, when considered alongside the relatively high branching fraction, indicates that this channel is a valuable addition to the suite of D decays that can be harnessed for the measurement of the unitarity-triangle angle γ through the process $B^\mp \rightarrow DK^\pm$. The decays $D \rightarrow K_{S,L}^0 \pi^+ \pi^-$ have also been employed as a tag to measure the CP contents of the modes $D \rightarrow \pi^+ \pi^- \pi^0$ and $D \rightarrow K^+ K^- \pi^0$. The results confirm the conclusion of a previous analysis [1], based on CP-eigenstate tags, and also suggested by earlier amplitude-model studies [24–26], that the CP-even content of the $\pi^+ \pi^- \pi^0$ final state is very high, and therefore this

decay too is a powerful mode for the measurement of γ . Combining the two sets of measurements yields $F_+^{\pi\pi\pi^0} = 0.973 \pm 0.017$ and $F_+^{KK\pi^0} = 0.732 \pm 0.055$. Now that their CP-even fractions have been measured, all three decay modes may also be used for studies of indirect CP violation and mixing in the $D^0 \bar{D}^0$ system [5].

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References

- [1] M. Nayak, et al., First determination of the CP content of $D \rightarrow \pi^+ \pi^- \pi^0$ and $D \rightarrow K^+ K^- \pi^0$, Phys. Lett. B 740 (2015) 1, arXiv:1410.3964 [hep-ex].
- [2] K.A. Olive, et al., Particle Data Group, Review of particle physics, Chin. Phys. C 38 (2014) 090001.
- [3] R. Aaij, et al., LHCb Collaboration, Search for CP violation in $D^0 \rightarrow \pi^+ \pi^- \pi^0$ decays with the energy test, Phys. Lett. B 740 (2015) 158, arXiv:1410.4170 [hep-ex].
- [4] R. Aaij, et al., LHCb Collaboration, Model-independent search for CP violation in $D^0 \rightarrow K^- K^+ \pi^- \pi^+$ and $D^0 \rightarrow \pi^- \pi^+ \pi^- \pi^+$ decays, Phys. Lett. B 726 (2013) 623, arXiv:1308.3189 [hep-ex].
- [5] S. Malde, C. Thomas, G. Wilkinson, Measuring CP violation and mixing in charm with inclusive self-conjugate multibody decay modes, arXiv:1502.04560 [hep-ph].
- [6] A. Bondar, A. Poluektov, The use of quantum correlated D^0 decays for ϕ_3 measurement, Eur. Phys. J. C 55 (2008) 51, arXiv:0801.0840 [hep-ex].
- [7] M. Gronau, D. London, How to determine all the angles of the unitarity triangle from $B_q^0 \rightarrow DK_S^0$ and $B_S^0 \rightarrow D\phi$, Phys. Lett. B 253 (1991) 483; M. Gronau, D. Wyler, On determining a weak phase from CP asymmetries in charged B decays, Phys. Lett. B 265 (1991) 172.
- [8] Y. Kubota, et al., CLEO Collaboration, The CLEO II detector, Nucl. Instrum. Methods A 320 (1992) 66; D. Peterson, et al., The CLEO III drift chamber, Nucl. Instrum. Methods A 478 (2002) 142; M. Artuso, et al., Construction, pattern recognition and performance of the CLEO III LiF-TEA RICH detector, Nucl. Instrum. Methods A 502 (2003) 91; R.A. Briere, et al., CLEO-c/CESR-c Taskforces and CLEO-c Collaboration, CLEO-c and CESR-c: a new frontier of weak and strong interactions, Cornell LEPP Report CLNS Report No. 01/1742, 2001.
- [9] D.J. Lange, The EvtGen particle decay simulation package, Nucl. Instrum. Methods A 462 (2001) 152.
- [10] R. Brun, et al., GEANT 3.21, CERN Program Library Long Writeup W5013, unpublished.
- [11] N. Lowrey, et al., CLEO Collaboration, Determination of the $D^0 \rightarrow K^- \pi^+ \pi^0$ and $D^0 \rightarrow K^- \pi^+ \pi^+ \pi^-$ coherence factors and average strong-phase differences using quantum-correlated measurements, Phys. Rev. D 80 (2009) 031105(R), arXiv:0903.4853 [hep-ex].
- [12] D.M. Asner, et al., CLEO Collaboration, Determination of the $D^0 \rightarrow K^+ \pi^-$ relative strong phase using quantum-correlated measurements in $e^+ e^- \rightarrow D^0 \bar{D}^0$ at CLEO, Phys. Rev. D 78 (2008) 012001, arXiv:0802.2268 [hep-ex].
- [13] J. Libby, et al., New determination of the $D^0 \rightarrow K^- \pi^+ \pi^0$ and $D^0 \rightarrow K^- \pi^+ \pi^+ \pi^-$ coherence factors and average strong-phase differences, Phys. Lett. B 731 (2014) 197, arXiv:1401.1904 [hep-ex].
- [14] H. Albrecht, et al., ARGUS Collaboration, Search for hadronic $b \rightarrow u$ decays, Phys. Lett. B 241 (1990) 278.
- [15] T. Skwarnicki, A study of the radiative cascade transitions between the Υ and Υ' resonances, PhD thesis (Appendix E), Cracow, INP, 1986, DESY F31-86-02.
- [16] Q. He, et al., CLEO Collaboration, Comparison of $D \rightarrow K_S^0 \pi$ and $D \rightarrow K_L^0 \pi$ decay rates, Phys. Rev. Lett. 100 (2008) 091801, arXiv:0711.1463 [hep-ex].
- [17] Y. Amhis, et al., HFAG, Averages of b-hadron, c-hadron, and τ -lepton properties as of summer 2014, arXiv:1412.7515, online updates at <http://www.slac.stanford.edu/xorg/hfag>.
- [18] J. Libby, et al., CLEO Collaboration, Model-independent determination of the strong-phase difference between the decays D^0 and $\bar{D}^0 \rightarrow K_{S,L}^0 h^+ h^-$ ($h = \pi, K$) and its impact on the measurement of the CKM angle γ , Phys. Rev. D 82 (2010) 112006, arXiv:1010.2817 [hep-ex].
- [19] B. Aubert, et al., BaBar Collaboration, Improved measurement of the CKM angle γ in $B^\mp \rightarrow D^{(*)} K^{(*)\mp}$ decays with a Dalitz plot analysis of D decays to $K_S^0 \pi^+ \pi^-$ and $K_S^0 K^+ K^-$, Phys. Rev. D 78 (2008) 034023, arXiv:0804.2089 [hep-ex].

- [20] P. del Amo Sanchez, et al., BaBar Collaboration, Evidence for direct CP violation in the measurement of the CKM angle γ with $B^{\mp} \rightarrow D^{(*)}K^{(*)\mp}$ decays, Phys. Rev. Lett. 105 (2010) 121801, arXiv:1005.1096 [hep-ex].
- [21] A. Poluektov, et al., Belle Collaboration, Evidence for direct CP violation in the decay $B \rightarrow D^{(*)}K$, $D \rightarrow K_S^0\pi^+\pi^-$ and measurement of the CKM phase ϕ_3 , Phys. Rev. D 81 (2010) 112002, arXiv:1003.3360 [hep-ex].
- [22] B. Aubert, et al., BaBar Collaboration, Measurement of the Cabibbo–Kobayashi–Maskawa angle γ in $B^{\mp} \rightarrow D^{(*)}K^{\mp}$ decays with a Dalitz analysis of $D^0 \rightarrow K_S^0\pi^-\pi^+$, Phys. Rev. Lett. 95 (2005) 121802, arXiv:hep-ex/0504039.
- [23] S. Brisbane, CLEO-c $D \rightarrow K_{S/L}^0\pi^+\pi^-$ binned Dalitz-plot analyses optimised for CKM angle γ measurement and the commissioning of the LHCb front-end electronics, DPhil thesis (Chapter 3), University of Oxford, 2010, CERN-THESIS-2010-206.
- [24] D. Cronin-Hennessey, et al., CLEO Collaboration, Searches for CP violation and $\pi\pi$ S wave in the Dalitz-plot analysis of $D^0 \rightarrow \pi^+\pi^-\pi^0$, Phys. Rev. D 72 (2005) 031102(R), arXiv:hep-ex/0503052; D. Cronin-Hennessey, et al., CLEO Collaboration, Phys. Rev. D 75 (2007) 119904 (Erratum).
- [25] B. Aubert, et al., BaBar Collaboration, Measurement of CP violation parameters with a Dalitz plot analysis of $B^{\pm} \rightarrow D_{\pi^+\pi^-\pi^0}K^{\pm}$, Phys. Rev. Lett. 99 (2007) 251801, arXiv:hep-ex/0703037.
- [26] M. Gaspero, B. Meadows, K. Mishra, A. Soffer, Isospin analysis of D^0 decay to three pions, Phys. Rev. D 78 (2008) 014015, arXiv:0805.4050 [hep-ex].